Robust technique for spectroscopic plasma analysis with application in real-time arc welding quality monitoring

J. Mirapeix A. Cobo O. M. Conde Universidad de Cantabria Grupo de Ingeniería Fotónica Department TEISA Avda. Los Castros s/n 39005 Santander, Spain E-mail: mirapeixjm@unican.es

C. Jaúregui

University of Southampton Optoelectronics Research Centre SO17 1BJ, United Kingdom

J. M. López-Higuera, MEMBER SPIE

Universidad de Cantabria Grupo de Ingeniería Fotónica

1 Introduction

Department TEISA Avda. Los Castros s/n 39005 Santander, Spain

Quality assurance of the welding process is of major importance in some critical applications, such as those within the aerospace sector, in which extensive nondestructive testing (NDT) procedures are routinely performed to every welded part. For these applications, the early detection of welding defects by means of an on-line quality monitoring system could avoid expensive NDT of early-identified defective weld seams.

Several monitoring principles have been demonstrated for this particular process, such as the measurement of the charge voltage induced on the welding nozzle,¹ acoustic emissions,² or imaging in different spectral bands (including infrared thermography³). It is worth noting that some welding methods, like arc welding or laser welding in the so-called "keyhole" regime, exhibit an inherent plasma emission in the ultraviolet (UV) and visible bands. The analysis of the emitted light offers relevant information about the welding process itself and, ultimately, about the quality of the resulting weld.⁴

These optical approaches to quality monitoring exhibit some advantages over electrical or acoustic methods, one of the most important being their immunity to the strong electromagnetic interferences (EMI) typically found in these processes, especially if the light is collected far from the welding torch. This is usually not a problem due to the high intensity from the thermal plasma, and setups with remote photodiodes⁵ and optical fiber capturing⁶ have been

Abstract. A new real-time implementation of the Boltzmann-plot technique for optical spectroscopic analysis of thermal plasmas is presented. Applications are found in the monitoring of industrial processes in which plasma emission naturally occurs, such as arc and laser welding. Experimental results show the feasibility of using the Boltzmann-plot technique for real-time calculation of the electronic temperature of the plasma, a parameter that has been correlated with common defects found in the weld seam. The foundations of the proposed technique, its implementation, and its capability for the on-line detection of defects in tungsten inert gas (TIG) welding of steel parts are discussed. © *2006 Society of Photo-Optical Instrumentation Engineers.* [DOI: 10.1117/1.2336416]

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demonstrated. Even the focusing optics of a welding laser system can be used to back-propagate the plasma emission from the welding spot.⁷

Another advantage of the optical methods refers to the rich information they can offer of the welding process itself. In particular, the spectroscopic analysis of the plasma emission reveals the characteristic and complex pattern of multiple emission peaks due to the multiple atomic species (at different degrees of ionization) contributing to the plasma. The processing of the collected spectrum involves typically the estimation of the amplitude and width of those peaks, along with the identification of the associate atomic species. From this data, relevant information of the plasma composition can be derived, which in turn can be correlated with some properties of the resulting welding seam.⁸

The most widely employed technique when a plasma is spectroscopically analyzed is the calculation of its electronic temperature, as it has been demonstrated that there is a relationship between this parameter and the resulting quality of the weld seam.^{8–10} The reliable acquisition and subsequent analysis of the electronic temperature of the plasma T_e allows the identification of defects that can appear in the weld seam as a result of a wide variety of causes: intensity fluctuations, shielding gas flux variations, irregularities in the width of the welded materials, and many others, not to mention those defects generated by human operators.

Although the Boltzmann-plot method permits a precise calculation of the plasma electronic temperature,^{4,11} the simplification that estimates T_e from a single pair of emis-

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sion lines of a given species is usually assumed valid. Computational cost is the main constriction when this solution is adopted, especially if real-time operation is required, as it is desirable when a sensor system is designed not only to detect possible defects, but also to actuate over the process to assure a result with a specific quality. In a previous work,¹² a real-time monitoring principle

In a previous work,¹² a real-time monitoring principle for the analysis of thermal plasmas was proposed. With this technique, the central peak wavelength estimation and, hence the spectral line identification, was performed by using the lineal phase operator (LPO) subpixel technique.¹³ By substituting the traditional solutions used for the spectral lines characterization, like the Levenberg-Marquardt technique, this approach allows the real-time calculation of the electronic temperature of the plasma T_e using a single pair of spectral lines. While the use of the LPO algorithm allows a more accurate estimation of the central wavelength associated to each spectral line, a cubic-spline interpolation of the noisy data is performed to obtain the peak intensities.

In this work, we propose to extend the previously described technique to multiple spectral lines using the Boltzmann-plot method for the T_e calculation. The results show that a robust estimation of T_e is obtained, while still being able to operate in real time. The proposed technique has been applied to on-line defect detection in arc welding processes. The correlation between the calculated T_e and defects produced in the weld seams are analyzed, along with performance comparisons with the previously reported method based on single pairs of spectral lines. The experimental results of tungsten inert gas (TIG) welds performed on steel plates show that the proposed technique is able to detect, in real-time, common defects such as oxidation due to insufficient flux of shielding gas, or lack (or excess) of penetration of the weld seam.

2 Theoretical Background

The detection of defects in the weld seam by means of spectroscopic analysis implies the necessity of finding a correlation between the weld quality and its defects and the electronic temperature of the plasma. This parameter can be obtained from the relative intensities of some emission lines belonging to a particular atomic species within those participating in the plasma.⁴ If the electron density is high enough so that:

$$N_e \ge 1.6 \times 10^{12} T_e^{1/2} (\Delta E)^3, \tag{1}$$

(where N_e is the electronic density, T_e is the electronic temperature, and ΔE is the largest energy gap in the atomic energy level system), then it can be assumed that the plasma is in local thermal equilibrium (LTE), which implies that the various particles within the plasma have Maxwellian energy distributions and collisional processes exceed radiative ones.

When LTE is satisfied, the population of an excited level can be determined with the Boltzmann equation:

$$N_m = \frac{N}{Z} g_m \exp\left(\frac{-E_m}{\kappa T_e}\right),\tag{2}$$

where N is the population density of the state m, Z is the partition function, g_m is the statistical weight, E_m is the



Real-Time Electronic Temperature

Fig. 1 Experimental setup.

excitation energy, and κ is the Boltzmann constant.

In optically thin plasmas, the intensity of a given spectral line I_{mn} , induced by a transition from level *m* to level *n*, can be related to the population density of the upper level N_m through

$$I_{mn} = N_m A_{mn} h \gamma_{mn}, \tag{3}$$

with A_{mn} being the transition probability, and $h\gamma_m$ the energy of such transition.

Combining Eqs. (2) and (3), T_e can be obtained from the following expression,

$$\ln\left(\frac{I_{mn}\lambda_{mn}}{A_{mn}g_m}\right) = \ln\left(\frac{hcN}{Z}\right) - \frac{E_m}{kT_e},\tag{4}$$

as the plot resulting from using various lines and representing the left-hand side of Eq. (4) versus E_m has a slope inversely proportional to T_e .

A common simplification in the T_e calculation process is usually adopted by considering only a pair of spectral lines of a given species. Thus, by knowing the ratio of the relative intensities of those lines and using Eq. (4), T_e can be calculated by means of the following expression

$$T_{e} = \frac{E_{m}(2) - E_{m}(1)}{k \ln \left[\frac{I(1)A(2)g_{m}(2)\lambda(1)}{I(2)A(1)g_{m}(1)\lambda(2)} \right]}$$
(5)

From Eq. (5), it is clear that a precise estimation of the central peak wavelength of each considered emission line is needed, especially considering that this value will be used in the identification of the chemical species associated to that line. As was mentioned before, the utilization of the linear phase operator subpixel algorithm will allow an efficient fulfilment of this step.¹²

To check the real-time behavior and performance of this technique, several experiments in the laboratory have been performed. They have been based on TIG welding of steel plates, with common defects intentionally produced. The experimental setup and results are discussed in the next section.



Fig. 2 Spectral lines within the range to be considered.

3 Experimental Setup and Results

The experimental setup, designed around a TIG welding system, is schematically described in Fig. 1. The power source of the welding system is the Kemppi Mastertig 2200, which offers output intensities from 5 to 220 A. The motion of the weld system is obtained by keeping the torch fixed and displacing the plates to be welded with a high-precision (1 μ m in both axes) positioning system, formed by the Newport MM4005 controller and two Newport MTM100PP1 linear stages. Argon, with typical flows of 12 L/min, has been used as a shielding gas. 1-mm-diam tungsten rods and AISI-304 stainless steel plates have been used in our experiments, although the consideration of any other material does not imply the necessity of a previous calibration of the system.

A collimating lens and a solarization-resistant optical fiber (Ocean Optics P400-UV-SR) were respectively used to acquire the light radiation from the plasma and to guide it to the entrance of the spectrometer. A low-cost, 2048 linear charge-coupled device (CCD) spectrometer (Ocean Optics



Fig. 3 Example of a reference weld seam without defects.



Fig. 4 Detection of weld seam defects with different techniques.

USB2000), with 0.3 nm of spectral resolution and a spectral range from 195 to 535 nm, was chosen.

Although spectral lines from various species were identified, Ar II and Fe I lines have been used for the estimation of T_e . Particularly, Ar II emission lines located from 459 to 506 nm, and Fe I lines from 372 to 416 nm, have been systematically used in the Boltzmann-plot technique. In Fig. 2, a typical spectrum obtained from the plasma of an arc weld is presented.

Although the final objective of the proposed system is to detect defects in weld seams, it was considered necessary to previously obtain a reference weld without defects. In Fig. 3, an example of a weld seam performed with a constant intensity of 42 A and a constant gas flux of 12 L/min is shown. The T_e signal depicted in Fig. 3(b) was calculated with four Ar II lines (located at 461, 480, 485, and 488 nm) as input data for the Boltzmann-plot process. After an ini-



Fig. 5 Detection of weld seam crater.

tial transient state due to the instability at the beginning of the welding process, the T_e signal shows no relevant variations, as no visible defects are to be found.

In Fig. 4, a weld seam with a small defect is presented. This defect was caused on a plate with a thickness of 1 mm by a reduction in the gas flux from 12 L/min to approximately 6 L/min at $x \approx 5$ cm. The resulting weld seam is shown in Fig. 4(a). Figures 4(b) (Ar II lines located at 459 and 488 nm) and 4(c) (Ar II lines located at 459 and 485 nm) present the result of calculating T_e with the classical approach indicated in Eq. (5). The selection of those lines was performed by trying to maximize the difference between their excitation energies E_m , which leads to a more precise T_e determination.¹⁴ It can be observed in Figs. 4(b) and 4(c) that the defect created by the reduction in the gas flux in the weld seam is identified, but in the first case the change in the T_e profile is too small to be detected.

Additionally, the performance of the proposed technique has been checked by using the Boltzmann plot with six Ar II lines (located at 404, 459, 473, 476, 485, and 488 nm) to calculate T_e . As expected, the sensitivity to defects in-



FIG. 6 Detection of gas cut

creases, as can be observed in Fig. 4(d). It must be noted that not only the change in the T_e profile is more pronounced when using the Boltzmann plot, but also the T_e profile, when no defects are present, is clearly less noisy than the ones depicted in Figs. 4(b) and 4(c).

A different test is presented in Fig. 5. In Figs. 5(a) (top) and 5(b) (bottom), the resulting weld seam is shown. In this case, a plate with a thickness of 1 mm was used, and the reduction of the gas flux resulted in a crater in the weld seam at $x \approx 3.5$ cm. The welding intensity is also varied at x=2.5 cm from 44 to 34 A, whose influence in the weld seam is obvious in Fig. 5(b). Although this discontinuity in the weld seam is not clearly identified in Fig. 5(c), where T_e has been calculated with a pair of Ar II emission lines (located at 473 and 485 nm), the crater is detected. In Fig. 5(d), the Boltzmann plot with four Ar II lines (with wavelengths of 473, 476, 480, and 485 nm) is used to obtain T_e . In this case, both the variation of intensity and the crater can be distinguished, as the first peak at x=2.5 cm with

the increase in weld penetration due to the gas flux reduction, and the third one at x=3.5 cm with the beginning of the crater.

As a final test, a 3-mm-thick plate was used. The analysis was performed using a constant intensity of 45 A and reducing the shielding gas flux from 12 to 2 L/min at $x \approx 5.5$ cm. In Figs. 6(a) and 6(b), the resulting weld seam is presented, where the defect caused by the gas flux reduction is obvious. In Fig. 6(b), the previously mentioned incision on the plate is clearly observed. In Fig. 6(c), where the lines of Ar II with wavelengths of 461 and 475 have been used, the defect at $x \approx 5.5$ cm is clearly recognized. However, the T_e obtained in Fig. 6(d) by using the Boltzmann plot with four Ar II lines (located at 461, 473, 476, and 480 nm) allows an easier identification of the defects, as T_e is more stable when no defects are to be found.

The previously described results from the welding tests confirm that the real-time implementation of the Bolztmann-plot technique offers a robust and fast calculation of the electronic temperature of the plasma, and with better sensitivity regarding defect detection than methods based on single pairs of spectral lines.

Further research is being conducted at present to collect the light with an optical fiber embedded in the welding torch, thus allowing a noninvasive monitoring system more appropriate for a production environment. Field tests will be also presented and discussed.

4 Conclusions

A new real-time implementation of the Boltzmann-plot technique for optical spectroscopic analysis of thermal plasmas, with applications in on-line quality monitoring of arc welding processes, is presented. Although the Boltzmann-plot technique is a well-known method to calculate the electronic temperature of a plasma generated in an arc welding procedure, it is commonly not used in on-line quality assessment systems due to its high computational cost. The classical approach is to obtain T_e by using only a single pair of emission lines from the plasma, but by involving several different lines in the Boltzmann plot, a more precise estimation of T_e is expected to be obtained, and hence, an easier discrimination of the defects appearing along the weld seam. The experimental results, based on

TIG welding of steel plates, confirm not only the feasibility of using the Boltzmann plot for the T_e calculation process in a real-time system, but also that it offers an easier way to identify weld defects.

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