

# Evidence for the direct decay of the 125 GeV Higgs boson to fermions

The CMS Collaboration<sup>†</sup>

**The discovery of a new boson with a mass of approximately 125 GeV in 2012 at the Large Hadron Collider<sup>1–3</sup> has heralded a new era in understanding the nature of electroweak symmetry breaking and possibly completing the standard model of particle physics<sup>4–9</sup>. Since the first observation in decays to  $\gamma\gamma$ ,  $WW$  and  $ZZ$  boson pairs, an extensive set of measurements of the mass<sup>10,11</sup> and couplings to  $W$  and  $Z$  bosons<sup>11–13</sup>, as well as multiple tests of the spin-parity quantum numbers<sup>10,11,13,14</sup>, have revealed that the properties of the new boson are consistent with those of the long-sought agent responsible for electroweak symmetry breaking. An important open question is whether the new particle also couples to fermions, and in particular to down-type fermions, as the current measurements mainly constrain the couplings to the up-type top quark. Determination of the couplings to down-type fermions requires direct measurement of the corresponding Higgs boson decays, as recently reported by the Compact Muon Solenoid (CMS) experiment in the study of Higgs decays to bottom quarks<sup>15</sup> and  $\tau$  leptons<sup>16</sup>. Here, we report the combination of these two channels, which results in strong evidence for the direct coupling of the 125 GeV Higgs boson to down-type fermions, with an observed significance of 3.8 standard deviations, when 4.4 are expected.**

The CMS and ATLAS experiments at the Large Hadron Collider (LHC) have reported the discovery of a new boson<sup>1–3</sup> with a mass near 125 GeV and with production rates, decay rates and spin-parity quantum numbers<sup>10–14</sup> compatible with those expected for the standard model Higgs boson<sup>4–9</sup>. In the standard model, the Higgs boson is a spin-zero particle predicted to arise from the Higgs field which is responsible for electroweak symmetry breaking<sup>17,18</sup>. As such, the standard model Higgs boson couples directly to the  $W$  and  $Z$  bosons, and indirectly to photons. To date, significant signals have been reported in channels where the boson decays to either  $\gamma\gamma$ ,  $WW$ , or  $ZZ$  boson pairs<sup>11–13</sup>, as predicted by the theory. Overall, these results directly demonstrate that the new particle is intimately related to the mechanism of spontaneous electroweak symmetry breaking, whereby the  $W$  and  $Z$  bosons become massive, and thus it has been identified as a Higgs boson.

The standard model also predicts that the Higgs field couples to fermions through a Yukawa interaction, giving rise to the masses of quarks and leptons. The structure of the Yukawa interaction is such that the coupling strength between the standard model Higgs boson and a fermion is proportional to the mass of that fermion. As the masses of many quarks and leptons are sufficiently well known from experiment, it is possible within the standard model to accurately predict the Higgs boson decay rates to these fermions. The existence of such decays and the corresponding rates remain to be established and measured by experiment. Indirect evidence for the Higgs coupling to the top quark, an up-type quark and

the heaviest elementary particle known to date, is implied by an overall agreement of the gluon–gluon fusion production channel cross-section with the standard model prediction. However, the masses of down-type fermions may come about through different mechanisms in theories beyond the standard model<sup>19</sup>. Therefore, it is imperative to observe the direct decay of this new particle to down-type fermions to firmly establish its nature. As a consequence of the Yukawa interaction discussed above, the most abundant fermionic Higgs boson decays will be to third-generation quarks and leptons, namely the bottom quark and the  $\tau$  lepton, as the decay of a Higgs boson with a mass around 125 GeV to top quarks is kinematically not allowed. Therefore, the most promising experimental avenue to explore the direct coupling of the standard model Higgs boson to fermions is in the study of the decay to bottom quark–antiquark pairs (denoted as  $b\bar{b}$ ) as well as to tau lepton–antilepton pairs (denoted as  $\tau\tau$ ).

Recently, the CMS Collaboration reported on a search for the decays of the new boson to  $b\bar{b}$  quark pairs<sup>15</sup> as well as to  $\tau\tau$  lepton pairs<sup>16</sup> based on data collected in 2011 and 2012. In this Letter, we report on the combination of the results from the study of these two decays to down-type fermion–antifermion pairs, performed for the first time at the LHC.

The CMS apparatus comprises several detectors specialized in identifying different types of particles. These detectors are arranged inside and outside a superconducting solenoid of 6 m internal diameter that provides a magnetic field of 3.8 T. The detector electronics process collision information at a rate of up to 40 MHz, and decide whether or not the crossing of proton beam bunches that took place at the centre of the detector produced proton–proton collisions of sufficient interest. Through a layered decision system, from the 20 MHz of proton bunch crossings, fewer than 1 kHz are saved for further analysis. A detailed description of the full CMS apparatus can be found in ref. 20. Data were collected at the LHC which delivered proton–proton collisions at centre-of-mass energies of 7 TeV (2011) and 8 TeV (2012). A total integrated luminosity of 5.1 (4.9) fb<sup>−1</sup> and 18.9 (19.7) fb<sup>−1</sup>, for the  $H \rightarrow b\bar{b}$  ( $H \rightarrow \tau\tau$ ) decay, has been analysed at 7 and 8 TeV, respectively. Each proton bunch crossing gave rise to a large number of simultaneous proton–proton collisions, on average 9 in 2011 and 21 in 2012. Such a large number of overlapping collisions presented exceptional challenges in reconstructing the individual particles produced in the collision where an interesting interaction took place. Those challenges were successfully met thanks to the CMS tracking system, able to separate collision vertices as close as 0.5 mm along the beam direction.

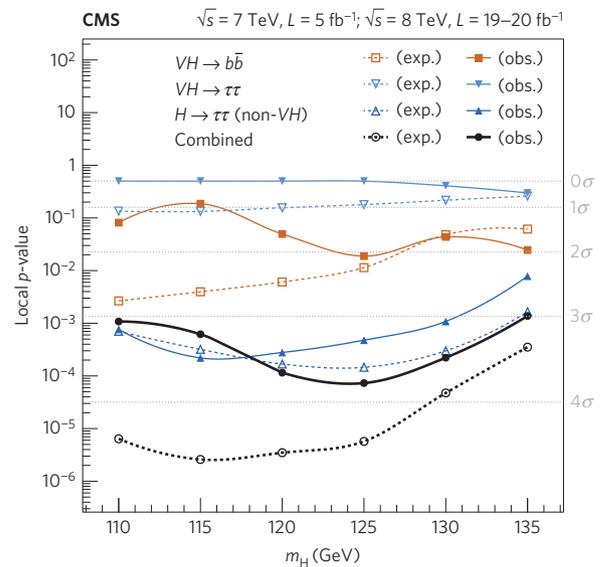
The Higgs boson lifetime is  $\sim 10^{-22}$  s. As a consequence, the detectors at the LHC only record the interactions of its decay products. For a Higgs boson mass of 125 GeV, the channels expected to be experimentally accessible include the decays to two photons, two  $W$  or  $Z$  bosons, a  $b\bar{b}$  quark pair and a  $\tau\tau$  lepton pair. The last two

<sup>†</sup>Full author list available online with Supplementary Information.

are especially challenging as the backgrounds due to standard model processes other than Higgs boson production are overwhelming and the invariant mass distributions of the reconstructed decay products do not produce sharp peaks. Techniques to estimate the background contributions directly from the collision data have been developed to face this challenge. The results are typically expressed in terms of upper limits on the product of the production and decay rates, significances for the background-only hypothesis to describe excesses as large, or larger, than those present in the data, and signal strengths relative to the expectation for the production and decay of the standard model Higgs boson. Both the observations in data and the expectations for the standard model Higgs boson are quoted. The expected quantities are important measures of the sensitivity of a particular search which do not depend on statistical fluctuations affecting actual observations in the data. Thus, they can also be used to compare the sensitivity of different experiments.

The CMS search for a Higgs boson decaying into a  $b\bar{b}$  quark pair is reported in ref. 15. The inclusive production of Higgs bosons in this decay channel is extremely hard to exploit owing to the overwhelming direct production of  $b\bar{b}$  quark pairs, as predicted by quantum chromodynamics (QCD). For this reason, present searches target the so-called ‘ $VH$  production’ mode, in which the Higgs boson is produced in association with a heavy vector boson,  $V$ , namely a  $W$  or a  $Z$  boson, identified via its leptonic decay. Requiring the presence of additional leptons with high transverse momentum provides an effective handle to reduce the QCD background. The Higgs boson candidate is reconstructed from two jets of particles, each of which has been identified as having a high probability of arising from the hadronization of a bottom quark. The identification is based on properties of hadrons containing bottom quarks, such as their long lifetimes, which lead to displaced vertices, or their decay to muons, both producing distinctive signatures that the CMS detector is sensitive to. Owing to the 10% mass resolution of the identified pair of  $b$ -quark jets<sup>15</sup>, the standard model Higgs boson signal is expected to produce a broad enhancement over the background in the  $b\bar{b}$  invariant mass distribution. The result of this analysis shows an excess over the background expectation for Higgs boson mass hypotheses,  $m_{\text{H}}$ , in the range of 120–135 GeV, with an observed (expected for the standard model Higgs boson) significance of 2.1 (2.1) standard deviations,  $\sigma$ , at a Higgs boson mass of 125 GeV. Correspondingly, the strength of this excess relative to the expectation for the standard model Higgs boson is  $1.0 \pm 0.5$ , where the uncertainty includes the effect of both statistical and systematic sources of uncertainty. A result obtained by the CDF and D0 experiments at the Tevatron proton–antiproton collider, which operated at a centre-of-mass collision energy of 1.96 TeV, initially revealed an excess of  $b\bar{b}$  events consistent with the mass of the Higgs boson observed at the LHC, yielding an observed significance of  $2.8\sigma$  for  $m_{\text{H}} = 125$  GeV (ref. 21). In a subsequent analysis the observed significance decreased, with the CDF and D0 experiments reporting an excess  $1.59^{+0.69}_{-0.72}$  times the expectation for the standard model Higgs boson with  $m_{\text{H}} = 125$  GeV (ref. 22). At the LHC, in previous studies based solely on the 2011 data, the CMS and ATLAS Collaborations have reported 95% confidence level (CL) upper limits of 2.5–5.5 times the standard model expectation for a Higgs boson in the 110–130 GeV mass range<sup>23,24</sup>.

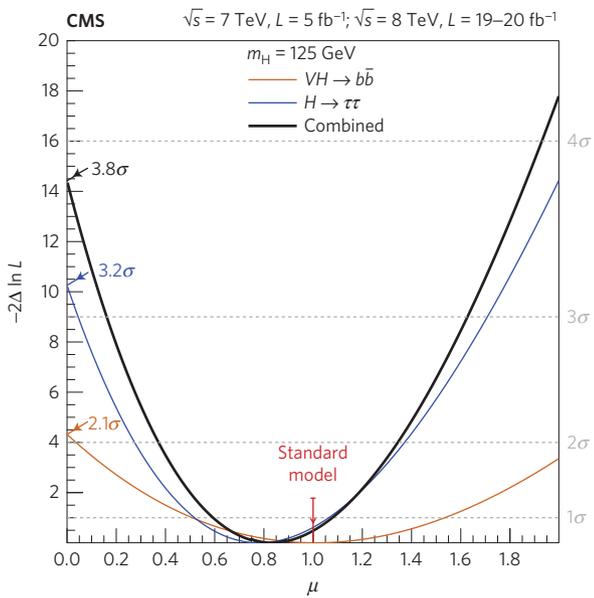
Evidence for a Higgs boson decaying to a  $\tau\tau$  lepton pair is reported by the CMS Collaboration in ref. 16. The lifetime of the  $\tau$  lepton is  $2.9 \times 10^{-13}$  s, short enough that  $\tau$  leptons decay before reaching any of the CMS detector systems. The predominant decay modes of  $\tau$  leptons are to a lighter lepton (that is, an electron or a muon) and two neutrinos, or to one or more pions and a neutrino<sup>25</sup>. In all cases, among the decay products there are neutrinos which cannot be directly detected. However, the presence of neutrinos among the products of an interaction can be inferred from an overall imbalance in the transverse momentum,



**Figure 1 | The  $p$ -value for the background-only hypothesis as a function of the Higgs boson mass hypothesis for both the observed data (obs.) and pseudo-data constructed using the median expectation for the standard model Higgs boson after the signal-plus-background fit to the data (exp.).** To help visualize the results, smooth lines connect the points at which the probabilities are evaluated. The  $H \rightarrow \tau\tau$  analysis results are broken down into a contribution from  $VH \rightarrow \tau\tau$  categories, which share the production process with the  $VH \rightarrow b\bar{b}$  analysis, and another contribution with all other  $H \rightarrow \tau\tau$  event categories. For every mass hypothesis tested, the non-fermionic decay contributions are taken to be those of the standard model Higgs boson with a mass of 125 GeV. Non-fermionic decays are considered as part of the background and consequently do not contribute to the measurement.

measured using the visible products of the interaction. The  $\tau$ -lepton decays involving pions are dominated by the decays to a single charged pion, a charged and a neutral pion, or three charged pions observed as a collimated jet of particles. All of these decay modes can be identified by the CMS detector, providing an important handle to suppress the QCD background. The main irreducible background is due to the copiously produced  $Z$  bosons which, like the Higgs boson, can decay to a  $\tau\tau$  lepton pair. The mass difference between the  $Z$  boson (91.2 GeV) and the region where a Higgs boson was observed ( $\sim 125$  GeV), together with a dedicated  $\tau\tau$  mass reconstruction technique<sup>16</sup>, allow a good separation of the two contributions. The gluon–gluon fusion production, the associated  $VH$  production, and the vector–boson fusion production processes are probed through the classification of the events according to kinematic and topological properties of the reconstructed particles. The result of the analysis shows a broad excess over background in the Higgs boson mass range of 115–130 GeV with an observed (expected) significance of  $3.2\sigma$  ( $3.7\sigma$ ) at a Higgs boson mass of 125 GeV. Correspondingly, the signal strength relative to the expectation for the standard model Higgs boson is found to be  $0.78 \pm 0.27$  at the Higgs boson mass of 125 GeV. Earlier, the CDF and D0 experiments at the Tevatron reported, for  $m_{\text{H}} = 125$  GeV, an observed (expected) 95% CL upper limit of 7.0 (5.7) times the standard model expectation<sup>22</sup>. Using the LHC data taken at 7 TeV, the ATLAS and CMS Collaborations reported, for  $m_{\text{H}} = 125$  GeV, observed (expected) 95% CL upper limits of 3.7 (3.5; ref. 26) and 4.2 (3.1; ref. 27) times the standard model expectation, respectively.

The CMS measurements in the  $VH \rightarrow b\bar{b}$  and  $H \rightarrow \tau\tau$  searches<sup>15,16</sup> are both consistent, within the precision of the present data, with each other and with the expectation for the production and decay of the standard model Higgs boson. Hence, as a next



**Figure 2 | Scan of the profile likelihood as a function of the signal strength relative to the expectation for the production and decay of a standard model Higgs boson,  $\mu$ , for  $m_H = 125$  GeV.** The statistical significance against the background-only ( $\mu = 0$ ) hypothesis is shown for the two channels and their combination. By definition, the expectation for the standard model Higgs boson with a mass of 125 GeV is  $\mu = 1$ . The non-fermionic decay contributions expected for the standard model Higgs boson with a mass of 125 GeV are considered as part of the background and therefore not scaled with  $\mu$ .

step, a combination of these two results, requiring the simultaneous analysis of the data selected by the two individual measurements, is performed. Such a simultaneous fit to both datasets accounts for all statistical and systematic uncertainties and their correlations. The statistical methods used for the combination were developed by the ATLAS and CMS Collaborations in the context of the LHC Higgs Combination Group. The presence of the signal is assessed by performing a fit to the observed data and using a test statistic based on the profile likelihood ratio. Systematic uncertainties are incorporated through nuisance parameters and are treated according to the frequentist paradigm. A detailed description of the methodology can be found in refs 28,29 and results presented here are obtained using asymptotic formulae<sup>30</sup>. To quantify an excess of events over the background expectation, the test statistic is defined as

$$q_0 = -2 \ln \frac{\mathcal{L}(\text{data} | b, \hat{\theta}_0)}{\mathcal{L}(\text{data} | \hat{\mu} \cdot s + b, \hat{\theta})} \quad \text{at } \hat{\mu} = 0$$

where the likelihood,  $\mathcal{L}$ , appearing in the numerator is that of the background-only hypothesis,  $s$  stands for the signal expected under the standard model Higgs boson hypothesis,  $\mu$  is a signal strength modifier introduced to accommodate deviations from standard model Higgs boson predictions,  $b$  stands for background contributions, and  $\theta$  are nuisance parameters describing the systematic uncertainties. The value  $\hat{\theta}_0$  maximizes the likelihood in the numerator under the background-only ( $\mu = 0$ ) hypothesis, while  $\hat{\mu}$  and  $\hat{\theta}$  define the point at which the likelihood reaches its global maximum in the signal-plus-background fit. The post-fit model, obtained after the signal-plus-background fit to the data, is used when deriving expected quantities. The post-fit model corresponds to the parametric bootstrap described in the statistics literature, and includes information gained in the fit regarding the values of all parameters<sup>31,32</sup>. The  $p$ -value for the background-only hypothesis is

**Table 1 | Summary of results for the Higgs boson mass hypothesis of 125 GeV.**

Channel ( $m_H = 125$ GeV)	Significance ( $\sigma$ )		Best-fit $\mu$
	Expected	Observed	
$VH \rightarrow b\bar{b}$	2.3	2.1	$1.0 \pm 0.5$
$H \rightarrow \tau\tau$	3.7	3.2	$0.78 \pm 0.27$
Combined	4.4	3.8	$0.83 \pm 0.24$

The  $p$ -values for the background-only hypothesis are expressed in terms of one-sided Gaussian tail significances and are provided in units of standard deviation ( $\sigma$ ). The expected significance is that obtained after the fit of the signal-plus-background hypothesis to the data. Note that the expected significance of  $2.1\sigma$  quoted in ref. 15 for the  $VH \rightarrow b\bar{b}$  channel was obtained before the fit of the signal-plus-background hypothesis to the data. The best-fit value of the signal strength relative to the expectation from the standard model,  $\mu$ , summarizes the profile likelihood scan of Fig. 2. For simplicity, uncertainties have been symmetrized. The statistical component represents more than 80% of the uncertainties.

calculated as the tail probability  $p_0 = P(q_0 \geq q_0^{\text{obs.}} | b)$ , where  $q_0^{\text{obs.}}$  is the value of the test statistic observed in data.

Whereas the expected Higgs boson signal in the  $VH \rightarrow b\bar{b}$  analysis consists exclusively of Higgs bosons decaying to down-type fermions, that is not the case in some event categories in the  $H \rightarrow \tau\tau$  analysis where a sizeable contribution from  $H \rightarrow WW$  decays is expected. The contribution from  $H \rightarrow WW$  decays has kinematic properties sufficiently different from those of  $H \rightarrow \tau\tau$  decays such that the two contributions do not overlap. To exclusively assess the fermionic decays of a Higgs boson, the expected contribution from  $H \rightarrow WW$  decays is considered as a background process. Given the discovery of a Higgs boson with properties as expected in the standard model and a mass close to 125 GeV, the  $pp \rightarrow H \rightarrow WW$  background contribution is taken from the expectation for the standard model Higgs boson with a mass of 125 GeV, including all associated uncertainties.

Figure 1 shows, as a function of the mass hypothesis for the Higgs boson decaying to fermions, the median expected and the observed  $p$ -value for the background-only hypothesis. The expectation is calculated after having performed the signal-plus-background fit to the observed data. The background-only hypothesis includes the  $pp \rightarrow H(125 \text{ GeV}) \rightarrow WW$  process for every value of  $m_H$ . The  $p$ -value can be expressed in terms of Gaussian tail probabilities and given in units of standard deviation ( $\sigma$ ), shown in the right vertical axis of Fig. 1. For all  $m_H$  values tested, the evidence against the background-only hypothesis is found to be  $3\sigma$  or more, with a maximum of  $3.8\sigma$  for  $m_H = 125$  GeV, corresponding to a  $p$ -value of  $7.3 \times 10^{-5}$ .

The study of the fermionic-decay contribution under the  $m_H = 125$  GeV hypothesis is particularly relevant because of the discovery of a Higgs boson with a mass near 125 GeV in the analysis of non-fermionic decays. As the mass resolution in the  $H \rightarrow b\bar{b}$  and  $H \rightarrow \tau\tau$  channels is, at best, around 10%, and the expected yields do not change significantly around  $m_H = 125$  GeV, the exact choice of  $m_H$  does not affect the conclusions. A scan of the profile likelihood is performed to estimate the best-fit signal strength relative to the standard model expectation, as well as confidence intervals. As described earlier, the  $H \rightarrow WW$  contribution is not treated as part of the signal. The result is shown in Fig. 2, from which a combined signal strength of  $0.83 \pm 0.24$  with respect to the standard model expectation is inferred for the production of a Higgs boson decaying to down-type fermions. The statistical component of the uncertainty represents 80% of the total uncertainty. This result is compatible with the expectation for the standard model Higgs boson, as summarized in Table 1.

In conclusion, the existing CMS searches for a Higgs boson decaying into bottom quarks and  $\tau$  leptons are consistent with the standard model prediction of a Yukawa structure, where

the fermionic couplings are proportional to the fermion masses. Combining the results of the CMS measurements in the  $H \rightarrow b\bar{b}$  and  $H \rightarrow \tau\tau$  decay channels and assuming that the Higgs boson is produced as expected in the standard model, reveals, for the first time, strong evidence for the direct coupling of the 125 GeV Higgs boson to down-type fermions, with an observed (expected) significance of 3.8 (4.4) standard deviations.

Received 25 January 2014; accepted 22 May 2014;  
published online 22 June 2014; corrected online 25 June 2014

## References

1. CMS Collaboration. Observation of a new boson at a mass of 125 GeV with the CMS experiment at the LHC. *Phys. Lett. B* **716**, 30–61 (2012).
2. CMS Collaboration. Observation of a new boson with mass near 125 GeV in pp collisions at  $\sqrt{s} = 7$  and 8 TeV. *J. High Energy Phys.* **06**, 081 (2013).
3. ATLAS Collaboration. Observation of a new particle in the search for the standard model Higgs boson with the ATLAS detector at the LHC. *Phys. Lett. B* **716**, 1–29 (2012).
4. Englert, F. & Brout, R. Broken symmetry and the mass of gauge vector mesons. *Phys. Rev. Lett.* **13**, 321–323 (1964).
5. Higgs, P. W. Broken symmetries, massless particles and gauge fields. *Phys. Lett.* **12**, 132–133 (1964).
6. Higgs, P. W. Broken symmetries and the masses of gauge bosons. *Phys. Rev. Lett.* **13**, 508–509 (1964).
7. Guralnik, G. S., Hagen, C. R. & Kibble, T. W. B. Global conservation laws and massless particles. *Phys. Rev. Lett.* **13**, 585–587 (1964).
8. Higgs, P. W. Spontaneous symmetry breakdown without massless bosons. *Phys. Rev.* **145**, 1156–1163 (1966).
9. Kibble, T. W. B. Symmetry breaking in non-Abelian gauge theories. *Phys. Rev.* **155**, 1554–1561 (1967).
10. CMS Collaboration. Study of the mass and spin-parity of the Higgs boson candidate via its decays to Z boson pairs. *Phys. Rev. Lett.* **110**, 081803 (2013).
11. CMS Collaboration. Measurement of the properties of a Higgs boson in the four-lepton final state. *Phys. Rev. D* **89**, 092007 (2014).
12. ATLAS Collaboration. Measurements of Higgs boson production and couplings in diboson final states with the ATLAS detector at the LHC. *Phys. Lett. B* **726**, 88–119 (2013).
13. CMS Collaboration. Measurement of Higgs boson production and properties in the WW decay channel with leptonic final states. *J. High Energy Phys.* **01**, 096 (2014).
14. ATLAS Collaboration. Evidence for the spin-0 nature of the Higgs boson using ATLAS data. *Phys. Lett. B* **726**, 120–144 (2013).
15. CMS Collaboration. Search for the standard model Higgs boson produced in association with a W or a Z boson and decaying to bottom quarks. *Phys. Rev. D* **89**, 012003 (2014).
16. CMS Collaboration. Evidence for the 125 GeV Higgs boson decaying to a pair of  $\tau$  leptons. Preprint at <http://arXiv.org/abs/1401.5041> (submitted, 2014).
17. Glashow, S. L. Partial-symmetries of weak interactions. *Nucl. Phys.* **22**, 579–588 (1961).
18. Weinberg, S. A model of leptons. *Phys. Rev. Lett.* **19**, 1264–1266 (1967).
19. Particle Data Group, Beringer J. *et al.* Review of particle physics. *Phys. Rev. D* **86**, 010001 (2012).
20. CMS Collaboration. The CMS experiment at the CERN LHC. *JINST* **3**, S08004 (2008).
21. CDF and D0 Collaborations. Evidence for a particle produced in association with weak bosons and decaying to a bottom–antibottom quark pair in Higgs boson searches at the Tevatron. *Phys. Rev. Lett.* **109**, 071804 (2012).
22. CDF and D0 Collaborations. Higgs boson studies at the Tevatron. *Phys. Rev. D* **88**, 052014 (2013).
23. CMS Collaboration. Search for the standard model Higgs boson decaying to bottom quarks in pp collisions at  $\sqrt{s} = 7$  TeV. *Phys. Lett. B* **710**, 284–306 (2012).
24. ATLAS Collaboration. Search for the standard model Higgs boson produced in association with a vector boson and decaying to a b-quark pair with the ATLAS detector. *Phys. Lett. B* **718**, 369–390 (2012).
25. Behringer, J. *et al.* (Particle Data Group) Status of Higgs Boson Physics. *Phys. Rev. D* **86**, 010001 (2012) and 2013 update for the 2014 edition; <http://pdg.lbl.gov>
26. ATLAS Collaboration. Search for the standard model Higgs boson in the  $H \rightarrow \tau^+\tau^-$  decay mode in  $\sqrt{s} = 7$  TeV pp collisions with ATLAS. *J. High Energy Phys.* **09**, 070 (2012).
27. CMS Collaboration. Search for neutral Higgs bosons decaying to  $\tau$  pairs in pp collisions at  $\sqrt{s} = 7$  TeV. *Phys. Lett. B* **713**, 68–90 (2012).
28. ATLAS and CMS Collaborations & LHC Higgs Combination Group. *Procedure for the LHC Higgs Boson Search Combination in Summer 2011*. Technical Report ATL-PHYS-PUB 2011-11, CMS NOTE 2011/005 (2011).
29. CMS Collaboration. Combined results of searches for the standard model Higgs boson in pp collisions at  $\sqrt{s} = 7$  TeV. *Phys. Lett. B* **710**, 26–48 (2012).
30. Cowan, G., Cranmer, K., Gross, E. & Vitells, O. Asymptotic formulae for likelihood-based tests of new physics. *Eur. Phys. J. C* **71**, 1554 (2011).
31. Efron, B. Bootstrap methods: Another look at the jackknife. *Ann. Statist.* **7**, 1–26 (1979) (See Section 8.k).
32. Lee, S. M. S. & Young, G. A. Parametric bootstrapping with nuisance parameters. *Stat. Probab. Lett.* **71**, 143–153 (2005).

## Acknowledgements

This Letter is dedicated to our dear colleague and friend Lorenzo Foà who passed away on 13 January 2014. Lorenzo's contribution to CMS was unique and left an indelible impact on the experiment in all its phases: Lorenzo was the CMS principal referee in the CERN LHC Committee during the approval process, initiated the participation of many groups during the formation of the CMS Collaboration, provided prescient advice and guidance during his period as CERN Director of Research, was Chair of the CMS Collaboration Board during the critical phase of construction, and was the first Chair of the Thesis Award Committee, enabling him to stay close to young physicists. We congratulate our colleagues in the CERN accelerator departments for the excellent performance of the LHC and thank the technical and administrative staffs at CERN and at other CMS institutes for their contributions to the success of the CMS effort. In addition, we gratefully acknowledge the computing centres and personnel of the Worldwide LHC Computing Grid for delivering so effectively the computing infrastructure essential to our analyses. Finally, we acknowledge the enduring support for the construction and operation of the LHC and the CMS detector provided by the following funding agencies: BMWF and FWF (Austria); FNRS and FWO (Belgium); CNPq, CAPES, FAPERJ and FAPESP (Brazil); MES (Bulgaria); CERN; CAS, MoST and NSFC (China); COLCIENCIAS (Colombia); MSES and CSF (Croatia); RPF (Cyprus); MoER, SF0690030s09 and ERDF (Estonia); Academy of Finland, MEC and HIP (Finland); CEA and CNRS/IN2P3 (France); BMBF, DFG and HGF (Germany); GSRT (Greece); OTKA and NIH (Hungary); DAE and DST (India); IPM (Iran); SFI (Ireland); INFN (Italy); NRF and WCU (Republic of Korea); LAS (Lithuania); MOE and UM (Malaysia); CINVESTAV, CONACYT, SEP and UASLP-FAI (Mexico); MBIE (New Zealand); PAEC (Pakistan); MSHE and NSC (Poland); FCT (Portugal); JINR (Dubna); MON, RosAtom, RAS and RFBR (Russia); MESTD (Serbia); SEIDI and CPAN (Spain); Swiss Funding Agencies (Switzerland); NSC (Taipei); ThEPCenter, IPST, STAR and NSTDA (Thailand); TUBITAK and TAEC (Turkey); NASU (Ukraine); STFC (United Kingdom); DOE and NSF (USA).

## Additional information

Supplementary information is available in the online version of the paper. Reprints and permissions information is available online at [www.nature.com/reprints](http://www.nature.com/reprints).

## Competing financial interests

The authors declare no competing financial interests.



This work is licensed under a Creative Commons

Attribution-NonCommercial-ShareAlike 3.0 Unported licence.

The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in the credit line; if the material is not included under the Creative Commons licence, users will need to obtain permission from the licence holder to reproduce the material. To view a copy of this licence, visit <http://creativecommons.org/licenses/by-nc-sa/3.0/>.